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Introduction to Geomagnetically Trapped Radiation

Martin Walt

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Chapter

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Appendix A

Summary of frequently used formulas

Numbers indicate where formula is located in text.

Motion of a charged particle in electric and magnetic fields

$$\mathbf{F} = q(\mathbf{v} \times \mathbf{B} + \mathbf{E}) \quad (2.1)$$

Gyroradius:

$$\rho = \frac{p_{\perp}}{Bq} = \frac{\gamma m_0 v_{\perp}}{Bq} \quad (2.4)$$

Gyrofrequency:

$$\begin{aligned} \Omega &= \frac{qB}{\gamma m_0} \\ \Omega \text{ (radians s}^{-1}\text{)} &= 1.758 \times 10^{11} \frac{B(T)}{\gamma} \quad \text{(electrons)} \\ &= 9.581 \times 10^7 \frac{B(T)}{\gamma} \quad \text{(protons)} \end{aligned}$$

Guiding center drift velocities

Electric field drift:

$$\mathbf{V}_E = \frac{\mathbf{E} \times \mathbf{B}}{B^2} \quad (2.8)$$

Gradient B drift:

$$\mathbf{V}_G = \frac{mv_{\perp}^2}{2qB^3} (\mathbf{B} \times \nabla B) \quad (2.23)$$

Curvature B drift:

$$\mathbf{V}_c = \frac{mv_{\parallel}^2}{qR_c} \frac{\mathbf{n} \times \mathbf{B}}{B^2} \quad (2.25)$$

In regions where $\nabla \times \mathbf{B} = 0$ the \mathbf{V}_c is given by

$$\mathbf{V}_c = \frac{mv_{\parallel}^2}{qB^3}(\mathbf{B} \times \nabla B) \quad (2.27)$$

Mirroring force:

$$\mathbf{F}_z = -\frac{mv_{\perp}^2}{2B} \frac{\partial B}{\partial z} \hat{\mathbf{e}}_z \quad (2.31)$$

$$= -\frac{\mu}{\gamma} \frac{\partial B}{\partial z} \hat{\mathbf{e}}_z \quad (4.20)$$

Dipole field equations ($B_0 = 3.12 \times 10^{-5}$ T)

$$B_r = -2B_0 \left(\frac{R_E}{r} \right)^3 \cos \theta \quad (3.13)$$

$$B_{\theta} = -B_0 \left(\frac{R_E}{r} \right)^3 \sin \theta \quad (3.14)$$

$$|B| = B_0 \left(\frac{R_E}{r} \right)^3 \sqrt{1 + 3 \cos^2 \theta} \quad (3.15)$$

Equation for dipole field line:

$$\begin{aligned} r &= R_0 \sin^2 \theta \\ &= R_0 \cos^2 \lambda \end{aligned} \quad (3.17)$$

Particle motion in geomagnetic field

Bounce period:

$$\tau_b(\text{s}) = 0.117 \left(\frac{R_0}{R_E} \right) \frac{1}{\beta} [1 - 0.4635(\sin \alpha_{\text{eq}})^{3/4}] \quad (4.28)$$

Drift period:

$$\tau_d(\text{s}) = C_d \left(\frac{R_E}{R_0} \right) \frac{1}{\gamma \beta^2} [1 - 0.333(\sin \alpha_{\text{eq}})^{0.62}] \quad (4.47)$$

where

$$\begin{aligned} C_d &= 1.557 \times 10^4 \text{ for electrons} \\ &= 8.481 \text{ for protons} \end{aligned}$$

Adiabatic invariants

$$J_1 = \mu = \frac{p_{\perp}^2}{2m_0 B} \quad (4.13)$$

$$J_2 = J = \oint \mathbf{p} \cdot d\mathbf{s} \quad (4.31)$$

$$J_3 = \Phi = q \int_S \mathbf{B} \cdot d\mathbf{S} \quad (4.50)$$

Relationship of flux $j(\alpha, E)$ to phase space density $F(\mathbf{q}, \mathbf{p})$

$$F(\mathbf{q}, \mathbf{p}) = j(\alpha, E)/p^2$$

Electric and magnetic field transformations

Lorentz transformations of electric and magnetic fields are given for a primed system moving at velocity \mathbf{V} with respect to an unprimed system. Parallel quantities are measured parallel to \mathbf{V} :

$$\mathbf{B}'_{\parallel} = B_{\parallel} \quad E'_{\parallel} = E_{\parallel}$$

$$\mathbf{B}'_{\perp} = \gamma \left(\mathbf{B}_{\perp} - \frac{\mathbf{V} \times \mathbf{E}}{c^2} \right)_{\perp} \quad \mathbf{E}'_{\perp} = \gamma (\mathbf{E}_{\perp} + \mathbf{V} \times \mathbf{B})_{\perp}$$

where

$$\gamma = (1 - V^2/c^2)^{-1/2}$$

In magnetospheric applications $V < 10^5 \text{ m s}^{-1}$, $10^{-8} \text{ T} < \mathbf{B} < 10^{-5} \text{ T}$, $\mathbf{E} < 0.1 \text{ vm}^{-1}$. Therefore, to an excellent approximation $\gamma = 1$, $B \gg \mathbf{V} \times \mathbf{E}/c^2$ and it is permissible to use

$$\mathbf{B}'_{\parallel} = B_{\parallel} \quad E'_{\parallel} = E_{\parallel}$$

$$\mathbf{B}'_{\perp} = \mathbf{B}_{\perp} \quad \mathbf{E}'_{\perp} = \mathbf{E}_{\perp} + (\mathbf{V} \times \mathbf{B})$$

Maxwell's equations

$$\nabla \cdot \mathbf{B} = 0 \quad \nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} - \frac{\nabla \cdot \boldsymbol{\rho}}{\epsilon_0}$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{i} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \quad \nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t}$$

where \mathbf{i} is the current density, ρ is charge density and $\boldsymbol{\rho}$ is the polarization of the medium.

Vector identities

\mathbf{A} , \mathbf{B} and \mathbf{C} are vector fields; $d\mathbf{S}$ is a surface element of S ; $d\mathbf{l}$ is a line element of a contour around S ; dV is an element of volume V :

$$(\mathbf{A} \times \mathbf{B}) \times \mathbf{C} = (\mathbf{A} \cdot \mathbf{C})\mathbf{B} - (\mathbf{B} \cdot \mathbf{C})\mathbf{A}$$

$$\nabla \times (\nabla \psi) = 0 \quad \nabla \cdot (\nabla \times \mathbf{F}) = 0$$

$$\iint_S (\nabla \times \mathbf{A}) \cdot d\mathbf{S} = \oint_C \mathbf{A} \cdot d\mathbf{l}$$

$$\iiint_V (\nabla \cdot \mathbf{A}) dV = \iint_S \mathbf{A} \cdot d\mathbf{S}$$

$$\mathbf{A} \cdot \nabla \mathbf{B} = (\mathbf{A} \cdot \nabla) \mathbf{B} = \left(A_x \frac{\partial}{\partial x} + A_y \frac{\partial}{\partial y} + A_z \frac{\partial}{\partial z} \right) \mathbf{B}$$

In spherical polar coordinates

$$\nabla \cdot \psi = \hat{\mathbf{e}}_r \frac{\partial \psi}{\partial r} + \hat{\mathbf{e}}_\theta \frac{1}{r} \frac{\partial \psi}{\partial \theta} + \hat{\mathbf{e}}_\phi \frac{1}{r \sin \theta} \frac{\partial \psi}{\partial \phi}$$

$$\nabla^2 \psi = \frac{1}{r^2} \left[\frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \right]$$

Relativistic relations

\mathbf{v} = velocity, Tm_0c^2 = kinetic energy, W = kinetic energy + rest mass energy

$$\beta = \frac{v}{c} \quad \gamma = (1 - \beta^2)^{-1/2} = T + 1$$

$$m = m_0\gamma \quad \mathbf{p} = m\mathbf{v} = m_0\gamma\mathbf{v}$$

$$W = (T + 1) m_0c^2 = (m_0^2c^4 + c^2p^2)^{1/2}$$

$$\beta^2 = \frac{T(T + 2)}{(T + 1)^2}$$