

# The Global Angular Form, The Euler Class, and the Thom Class

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**Ekin Kaan<sup>a</sup>**

*<sup>a</sup>Department of Mathematics, University of Vienna  
Universitatring 1, 1010 Wien*

ABSTRACT: In this paper, we will construct an explicit formula of Thom class for an arbitrary  $\mathbb{R}^2$ -bundle in local coordinates.

References:

Differential forms in Algebraic Topology - Bott and Tu (p 70-75)

Characteristic Classes - Milnor and Stasheff

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## 1 Recall: Definition of the Thom Class

Let  $M$  be an  $m$  dimensional manifold and  $\pi : E \rightarrow M$  be an  $\mathbb{R}^n$  oriented vector bundle. The Thom class  $\Phi$  of this vector bundle is defined as the image of 1 in  $H^0(M)$  under the Thom isomorphism

$$\mathcal{F} : H^k(M) \rightarrow H_{cv}^{k+n}(E).$$

**Proposition 1.** *The Thom class  $\Phi$  of the vector bundle given above is the unique cohomology class on  $H_{cv}^n(E)$  whose restriction to each fiber  $F$  is the generator of  $H_c^n(F)$ .*

This means that  $\Phi|_F$  is a bump form on the fiber with total integral 1.

## 2 The Global Angular Form, The Euler Class, and the Thom Class

\*Orientation: A top form on a manifold  $M$  is called positive if it is in the orientation class of  $M$ .

The deformation retraction  $\pi : \mathbb{R}^n - \{0\} \rightarrow S^{n-1}$  implies a group isomorphism

$$\pi^* : H^{n-1}(S^{n-1}) \rightarrow H^{n-1}(\mathbb{R}^n - \{0\}) \quad (2.1)$$

The image of the generator  $\sigma$  of  $H^{n-1}(S^{n-1})$ ,  $\psi = \pi^*\sigma$  under this isomorphism is called the angular form on  $\mathbb{R}^n - \{0\}$ . If  $\rho(r)$  is a continuously increasing function of the radius  $r$  which is  $-1$  near 0 and 0 near infinity,  $d\rho(r) = \frac{\partial \rho}{\partial r} dr$  is a bump function on  $\mathbb{R}$  with total integral 1. Therefore,  $d\rho \cdot \psi$  is a compactly supported form on  $\mathbb{R}^n$  and  $(d\rho \cdot \psi)$  is the generator of  $H_c^n(\mathbb{R}^n)$ . Since  $\psi$  is an element of the cohomology class,  $d\psi = 0$  and therefore,  $(d\rho \cdot \psi) = d(\rho \cdot \psi) = \Phi$  is the Thom class. To make sense of a function of  $r$ , we have to endow the total space with a Riemannian structure.

$$\langle . | . \rangle_E = \sum_{\alpha} \rho_{\alpha} \langle . | . \rangle_{\alpha} \quad (2.2)$$

Now we will generalize this construction to general  $\mathbb{R}^2$  bundle. Let  $E^0$  be the set of all nonzero elements of the total space  $E$  of an  $\mathbb{R}^2$  vector bundle over a manifold  $M$ . Locally,  $E^0|_{U_\alpha}$  has a frame  $\pi^*x_1, \dots, \pi^*x_m, r_\alpha, \theta_\alpha$  of  $U_\alpha \times \mathbb{R}^2$ . Note that we can always find a local frame when we go from one local subspace  $U_\alpha$  to  $U_\beta$  such that  $r_\alpha = r_\beta$  on the intersection  $U_\alpha \cap U_\beta$ . However, the angular coordinate  $\theta_\alpha$  will change with some angle depending on local trivialization of the vector bundle unless the vector bundle is trivial. Denote this change by an angle  $\pi^*\phi_{\alpha\beta}$  where  $0 \leq \phi_{\alpha\beta} \leq 2\pi$ .

$$\theta_\beta = \theta_\alpha + \pi^*\phi_{\alpha\beta} \quad (2.3)$$

\*\*\*  $\phi_{\alpha\beta}$  is defined on the manifold  $M$  instead of the total space  $E$ . Because it kind of depends on how curved the base space is.

Note that the  $\phi_{\alpha\beta}$ 's do not satisfy cocycle condition.

$$\phi_{\alpha\beta} + \phi_{\beta\gamma} - \phi_{\alpha\gamma} = 2\pi\epsilon_{\alpha\beta\gamma} \quad (2.4)$$

where  $\epsilon_{\alpha\beta\gamma} \in \mathbb{Z}$ .  $d\phi_{\alpha\beta}$  clearly satisfies cocycle condition.

Define a 1-form on  $U_\alpha$

$$\xi_\alpha = \frac{1}{2\pi} \sum_{\gamma} \rho_\gamma d\phi_{\alpha\gamma} \quad (2.5)$$

on open subset  $U_\alpha$  where  $\{\rho_\alpha\}$  is a partition of unity subordinate to the open cover  $\{U_\alpha\}$ .

On  $U_\alpha \cap U_\beta$

$$\xi_\beta - \xi_\alpha = \frac{1}{2\pi} \sum_{\gamma} \rho_\gamma d\phi_{\beta\gamma} - \frac{1}{2\pi} \sum_{\gamma} \rho_\gamma d\phi_{\alpha\gamma} \quad (2.6)$$

$$= \frac{1}{2\pi} \sum_{\gamma} \rho_\gamma (d\phi_{\beta\gamma} - d\phi_{\alpha\gamma}) \quad (2.7)$$

$$= \frac{1}{2\pi} \sum_{\gamma} \rho_\gamma d\phi_{\alpha\beta} = \frac{1}{2\pi} d\phi_{\alpha\beta} \quad (2.8)$$

It follows that  $d\xi_\alpha = d\xi_\beta$  on  $U_\alpha \cap U_\beta$ . It means that the local 2-forms which match on the intersection can be glued together to a global 2-form  $e$  on  $M$ . This global 2-form  $e \in H^2(M)$  is called the Euler class.

It is easy to check that the Euler class is independent of choice of 1 form  $\xi_\alpha$ .

Now we can define the global angular form on  $E^0$  which can be written in local coordinates as follows:

$$\psi|_{U_\alpha} = \frac{d\theta_\alpha}{2\pi} - \pi^*\xi_\alpha \quad (2.9)$$

The global angular form  $\psi$  is well-defined since

$$\frac{d\theta_\alpha}{2\pi} - \pi^*\xi_\alpha = \frac{1}{2\pi}d\theta_\beta - \frac{1}{2\pi}\pi^*\phi_{\alpha\beta} - \pi^*\xi_\beta + \frac{1}{2\pi}\pi^*d\phi_{\alpha\beta} \quad (2.10)$$

$$= \frac{d\theta_\beta}{2\pi} - \pi^*\xi_\beta \quad (2.11)$$

on the intersection  $U_\alpha \cap U_\beta$ . Note that the restriction of the global angular 1-form  $\psi$  to a fiber is the angular form  $\frac{1}{2\pi}d\theta$  and the global angular form is not closed:

$$d\psi = -\pi^*d\xi_\alpha = -\pi^*e \quad (2.12)$$

**Remark 1.** *If the vector bundle is trivial, then the Euler class vanishes and the global angular form is closed.*

The Euler class can be computed explicitly in terms of transition functions of the total space  $E$ . Let  $g_{\alpha\beta} : U_\alpha \cap U_\beta \rightarrow SO(2)$  be the transition functions. Since  $SO(2)$  can be identified with  $U(1)$ , we can think of  $g_{\alpha\beta}$  as the complex valued functions. Since an arbitrary element of  $U(1)$  can be written as  $e^{i\theta}$ , we have

$$\theta_\alpha - \theta_\beta = -\pi^*\phi_{\alpha\beta} = \pi^*\left(\frac{1}{i} \log g_{\alpha\beta}\right) \quad (2.13)$$

$$\pi^*d\phi_{\alpha\beta} = -\pi^*\left(\frac{1}{i}d \log g_{\alpha\beta}\right) \quad (2.14)$$

Since  $\pi$  is surjective,  $\pi^*$  is injective and therefore

$$d\phi_{\alpha\beta} = -\frac{1}{i}d \log g_{\alpha\beta} \quad (2.15)$$

Thus, the local 1-forms we defined before become

$$\xi_\alpha = \frac{1}{2\pi} \sum_{\gamma} \rho_\gamma d\phi_{\gamma\alpha} = -\frac{1}{2\pi i} \sum_{\gamma} \rho_\gamma d \log(g_{\gamma\alpha}) \quad (2.16)$$

Therefore, the Euler class becomes

$$e(E) = d\xi_\alpha = -\frac{1}{2\pi i} \sum_{\gamma} d(\rho_\gamma d \log g_{\gamma\alpha}) \quad (2.17)$$

on  $\pi^{-1}(U_\alpha)$ .

**Proposition 2.** *If  $f : N \rightarrow M$  is a  $C^\infty$  map between manifolds and  $E$  is an  $\mathbb{R}^2$  oriented bundle on  $M$ , then*

$$e(f^{-1}E) = f^*e(E) \quad (2.18)$$

**proof :** Since the transition of the  $f^{-1}(E)$  bundle is by definition  $f^*g_{\alpha\beta}$  and pullback commutes with exterior derivative, it follows from the equation 2.15.

Now we can similarly define a cohomology class and show that it is in fact the Thom class:

$$\Phi = d(\rho(r).\psi) = d\rho(r).\psi - \rho(r)\pi^*(e) \quad (2.19)$$

The fact that  $\psi$  is defined only on  $E^0$  does not matter as  $d\rho(r)$  vanishes as  $r$  approaches 0.

### Properties of $\Phi$

- $\Phi$  has a compact support in each fiber.
- $\Phi$  is closed:  $d\Phi = -d\rho.d\psi - d\rho.d\psi$
- restriction to each fiber has total integral 1

$$\int_0^\infty \int_0^{2\pi} d\rho(r) \frac{d\theta}{2\pi} = \rho(\infty) - \rho(0) = 1 \quad (2.20)$$

- the cohomology class of  $\Phi$  is independent of choice of  $\rho(r)$ . Suppose  $\rho'(r)$  is another function whose value is  $-1$  near 0 and 0 near infinity. Then

$$\Phi - \Phi' = d((\rho - \rho').\psi) \quad (2.21)$$

Since  $(\rho - \rho').\psi$  is a 1 form, it implies that the equivalence class of  $\Phi$  equals to that of  $\Phi'$ .

Therefore,  $\Phi$  defines the Thom class.

**Proposition 3.** *The pullback of the Thom class to  $M$  by zero section is the Euler class.*

**proof :**

Let  $s : M \rightarrow E$  be the zero section. Then

$$s^*\Phi = d(\rho(0)).s^*\psi - \rho(0)s^*\pi^*e = e \quad (2.22)$$

since  $s^*\pi^* = id_M$ .

Finally, we can write an explicit formula for the Thom class:

$$\Phi = d\rho.\psi - \rho d\psi \quad (2.23)$$

$$\psi = \frac{d\theta_\alpha}{2\pi} - \pi^*d\xi_\alpha \quad (2.24)$$

$$= \frac{d\theta_\alpha}{2\pi} + \frac{1}{2\pi i} \pi^* \sum_\gamma \rho_\gamma d \log g_{\alpha\beta} \quad (2.25)$$

Thus,

$$\Phi = d\left(\rho \frac{d\theta_\alpha}{2\pi}\right) + \frac{1}{2\pi i} d\rho \pi^* \sum_\gamma \rho_\gamma d \log g_{\alpha\beta} \quad (2.26)$$